Magnetic Manipulation of Lateral Migration Behavior of a Ferrofluid Droplet in a Plane Poiseuille Flow

Md Rifat Hassan¹, Cheng Wang^{1,*}

Department of Mechanical and Aerospace Engineering, Missouri S&T, Rolla, MO, USA *Corresponding author(wancheng@mst.edu)

INTRODUCTION: A detailed investigation on the effect of a uniform magnetic field on the lateral migration of a ferrofluid droplet in a plane Poiseuille flow by means of numerical simulation is presented here. In this case, the magnetic field is applied along different arbitrary directions.

UTATIONAL METHODS: The conservative level set

RESULTS: The effect of different arbitrary magnetic field directions on the final equilibrium position of the droplet



method is used to track the dynamic interface of the droplet where the level set function ϕ is advected by the velocity field[1,2]:

$$\frac{d\phi}{dt} + \mathbf{u} \cdot \nabla \phi = \gamma \nabla \cdot \left(\varepsilon \nabla \phi - \phi (1 - \phi) \frac{\nabla \phi}{|\nabla \phi|} \right)$$

Being treated as a single phase flow, the different properties of the flow domain are related to ϕ through the following equations:

$$\rho = \rho_c + (\rho_d - \rho_c)\phi; \quad \eta = \eta_c + (\eta_d - \eta_c)\phi$$
$$\mu = \mu_c + (\mu_d - \mu_c)\phi; \quad \chi = \chi_c + (\chi_d - \chi_c)\phi$$

The flow field under the effect of uniform magnetic field can be governed by the continuity and momentum equations:

$$\nabla \cdot \boldsymbol{u} = 0$$

$$\rho \left(\frac{\delta \mathbf{u}}{\delta t} + \mathbf{u} \cdot \nabla \mathbf{u} \right) = -\nabla p + \nabla \cdot \boldsymbol{\tau} + \mathbf{F}_{\sigma} + \mathbf{F}_{m}$$

Figure 2. Lateral migration of a ferrofluid droplet at $Re_d = 0.03$, $\lambda = 1$, and $Bo_m = 0$.

Figure 3. Effect of different magnetic field directions on the migration behavior of the droplet at $Re_d =$ $0.03, \lambda = 1$, and $Bo_m = 8.72$.



where, the surface tension force, \mathbf{F}_{σ} can be defined as:

 $\mathbf{F}_{\sigma} = \nabla \cdot [\sigma \{\mathbf{I} + (-\mathbf{n}\mathbf{n}^{T})\}\delta]$

and magnetic force, \mathbf{F}_m can be calculated as:

 $\mathbf{F}_m = \nabla \cdot \boldsymbol{\tau}_m = \nabla \cdot \left(\mu \mathbf{H} \mathbf{H}^T - \frac{\mu}{2} H^2 \mathbf{I} \right)$

The magneto-static Maxwell equation can be written as:

 $\nabla \cdot \mathbf{B} = 0, \ \nabla \times \mathbf{H} = 0, \ \nabla \cdot (\mu \nabla \varphi) = 0$ $\mathbf{M} = \chi \mathbf{H}$ and $\mathbf{B} = \mu_0 (\mathbf{H} + \mathbf{M}) = \mu_0 (1 + \chi) \mathbf{H}$







Figure 4. Steady state velocity, magnetic field profiles, and equilibrium droplet shapes at $Re_d = 0.03$, $\lambda = 1$, and $Bo_m = 8.72$.

CONCLUSIONS: In the absence of any external forces, at $\lambda = 1$, the droplet finds its equilibrium position at a location approximately 19 μ m below the center of the channel. Applying a magnetic field along arbitrary directions results in different equilibrium positions along

Figure 1. Schematic illustration of a ferrofluid droplet suspended in another medium in a Poiseuille flow under the application of a uniform magnetic field, **H**₀.

the channel due to disparate alignments of the droplet with the flow field. At $\alpha = 0^{\circ}$, the droplet is found to be closer to the bottom wall, while at $\alpha = 45^{\circ}$ the droplet settles closer to the center, and at $\alpha = 90^{\circ}$, the droplet finds its equilibrium position exactly at the center of the channel.

References

- COMSOL, "CFD Module Application Library Manual.".
- E. Olsson and G. Kreiss, "A conservative level set method for two 2. phase flow," J. Comput. Phys., vol. 210, no. 1, pp. 225–246, 2005.

Excerpt from the Proceedings of the 2019 COMSOL Conference in Boston